

ORTHOGONAL SYSTEMS OF FUNCTIONS (chapter 2)

We start from definition of scalar product (f, g)

$$(f, g) \stackrel{(*)}{=} \int_a^b f(x) \overline{g(x)} w(x) dx$$

There are several possible scalar products, depending on the choice of the interval $[a; b]$, and the weight function $w(x) \geq 0$
 (a very simple choice of weight is $w(x) = 1$ or $w(x) = c$ - a constant)

In classical Fourier Series $[a; b] = [-\pi; \pi]$, but it could be any interval of length $T > 0$ (period of f). The interval could also be $[0; +\infty)$ - a half-line or $(-\infty; +\infty)$ a whole line.

Chosen a specific scalar product

Def. we say that $f \perp g$ with respect to the scalar product $(*)$
 if and only if $\int_a^b f(x) \overline{g(x)} w(x) dx = 0$

(the integral = 0 because of "cancellation effect" of positive and negative areas)

Remark $(f, f) = \int_a^b |f(x)|^2 w(x) dx \geq 0$

" $\|f\|^2$ on $L^2([a, b], w)$

Consider a set of (∞ many) functions $\phi_n(x)$ with $x \in (a, b)$, and $n \in I$ (set of indices, usually $\mathbb{N}^+, \mathbb{N}^0, \mathbb{Z}$).

This is called orthogonal system if:

$$(1) \quad \phi_n \in L^2([a, b], w) \quad \forall n \in I$$

$$(2) \quad \|\phi_n\|_2 > 0 \quad \forall n \in I \quad \text{where}$$

(the norm must be positive)

$$\|\phi_n\|_2 = \|\phi_n\| = \left(\int_a^b |\phi_n(x)|^2 w(x) dx \right)^{\frac{1}{2}}$$

$$(3) \quad (\phi_n, \phi_m) = 0 \quad \text{if } n \neq m$$

(these functions are \perp to each other)

Example (important)

$$\phi_m(x) = e^{inx} = \cos mx + i \sin mx$$

$$I = \mathbb{Z} = \{-\dots, -3, -2, -1, 0, 1, 2, 3, \dots\}$$

$[a; b] = [-\pi; \pi]$ (actually any interval of length 2π is OK, e.g. $[0; 2\pi]$ or $[-\frac{\pi}{4}; \frac{7}{4}\pi]$)

$$w(x) = 1 \quad L^2([- \pi; \pi], dx)$$

$$(1) \int_{-\pi}^{\pi} |\phi_m(x)|^2 dx = \int_{-\pi}^{\pi} 1 dx = 2\pi \Rightarrow \|\phi_m\|_2 = \sqrt{2\pi}$$

$$|\phi_m(x)|^2 = |\cos mx + i \sin mx|^2 = \cos^2 mx + \sin^2 mx = 1$$

$$(2) \|\phi_m\|_2 = \sqrt{2\pi}$$

N.B.: (1) is almost trivial in this case, because ϕ_m itself is bounded and $\int_a^b |\phi_m(x)|^2 dx = \infty$ impossible, BUT for general orthogonal systems the $\phi_m(x)$ could be unbounded; as long as $\|\phi_m\| < \infty$ FINITE.

$$(3) (\phi_n, \phi_m) = \int_{-\pi}^{\pi} \underbrace{e^{inx}}_{\phi_n} \underbrace{e^{-imx}}_{\phi_m} dx = \int_{-\pi}^{\pi} e^{i(n-m)x} dx = \text{with } n-m \neq \text{(INTEGER)}$$

$$= \left[\frac{e^{i(n-m)x}}{i(n-m)} \right]_{-\pi}^{\pi} = 0 \quad \begin{matrix} \text{REMEMBER} \\ e^{i\pi} = e^{-i\pi} \end{matrix}$$

||| 0 because of cancellation

Def.: suppose $f(x) \in L^1([a, b], w)$ and also $f(x) \overline{\phi_m(x)} \in L^1([a, b], w)$ for $m \in I$, then

$$c_m = \frac{1}{\|\phi_m\|_2^2} \cdot (f, \phi_m)$$

are the generalized Fourier coefficients of $f(x)$ w.r.t. the orthogonal system $\{\phi_n\}$.

N.B. There are many orthogonal systems besides the classical e^{inx} ...

N.B. When $\phi_n(x) = e^{inx}$, $I = \mathbb{Z}$, then c_n are the classical Fourier coeff's of the 2π -periodic function f :

$$c_n = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-inx} dx$$

$\| \phi_n \|_2^2 \nearrow \underbrace{(f, e^{inx})}_{(f, g)}$

N.B. if instead of 2π -periodic functions we want to consider T -periodic functions, then $\phi_n(x) = e^{i \frac{2\pi}{T} nx}$, $[a, b] = [-\frac{T}{2}, \frac{T}{2}]$, or any interval of length T , like e.g. $[0, T]$.

In the T -periodic case, our scalar product will be: ($w=1$)

$$(f, g) = \int_{-T/2}^{T/2} f(x) \overline{g(x)} dx$$

$$c_n = \frac{1}{T} \int_{-T/2}^{T/2} f(x) e^{-i \frac{2\pi}{T} nx} dx$$

Remark F.S. are well adapted both to the case of T -periodic functions and the case of functions that are only defined in $[a; b]$, with ($b - a = T$), because we can always assume that a function defined only on $[a; b]$ is periodic outside.

Guideline (Road Map) to apply Fourier Theory to PDE's:

- (1) Domain has a side which is a finite interval \Rightarrow Fourier Series (classical or generalized)
- (2) Domain has a full line in the boundary \Rightarrow Fourier Transform
- (3) Domain has a half line on the boundary \Rightarrow Laplace Transform

The expression $\sum_{n \in I} c_n \phi_n(x)$ is called the expansion of $f(x)$

as Fourier Series (Generalized) w.r.t. the orthogonal system $\{\phi_n\}$.

This F.S. converges to $f(x)$ at least in the L^2 sense

e.g. $\lim_{N \rightarrow \infty} \left\| f - \sum_{n=-N}^N c_n \phi_n(x) \right\|_2 = 0$

Then, if f is "good enough" (for example $f \in C^1[a; b]$), then we have other kinds of convergence ($C^1 \Rightarrow$ pointwise, uniform...)

N.B. If $I = \mathbb{Z}$ the partial sums are from $-N$ to N (balanced, symmetric, finite sums). If $I = \{0, 1, 2, 3, \dots\}$, then the partial sums are from 0 to N , etc.

Def. An orthogonal system is called ortho-normal if

$$\|\phi_n\|_2 = 1 \text{ for } n \in I$$

Any orthogonal system can be normalized, just multiplying each ϕ_n by a suitable constant.

Theorem (minimizing property of Fourier Coefficients)

$$f(x) \sim \sum_{n \in I} c_n \phi_n(x) \quad f \in L^2([a; b]) \quad c_n = \frac{1}{\|\phi_n\|_2} \cdot (f, \phi_n)$$

$$\left\| f(x) - \sum_{j=-N}^N c_j \phi_j(x) \right\|_2 \leq \left\| f(x) - \sum_{j=-N}^N d_j \phi_j(x) \right\|_2$$

and the = holds $\Leftrightarrow d_j = c_j$ (Fourier coeff's $\forall j$)

We could say that the Fourier partial sums ($\sum_{-N}^N \dots$) give us the best L^2 projection on a vector sub-space of L^2 of finite dimension ($2N+1$) (Proof later...)

Corollary $\left\| f(x) - S_N(x) \right\|_2^2 = \|f\|_2^2 - \sum_N |c_j|^2 \|\phi_j\|_2^2$

Corollary of the Corollary (Bessel's inequality):

$$\sum_{j \in I} |c_j|^2 \|\phi_j\|_2^2 \leq \|f\|_2^2$$

because $\|f(x) - S_N(x)\|_2$ is positive, so $f(x) > S_N(x)$

Def. We say that $\{\phi_n\}$ orthogonal system has the Parseval Property if, for $\forall f \in L^2([a; b])$, we have the equality in the Bessel's inequality:

$$\sum_{j \in I} |c_j|^2 \|\phi_j\|_2^2 = \|f\|_2^2 \quad (*)$$

(∞ -dimensional Pythagorean theorem)

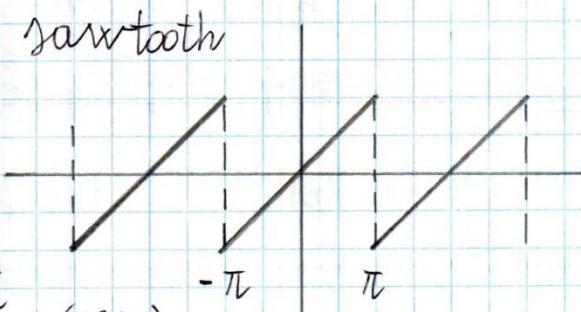
Def. An orthogonal system $\{\phi_n\}$ is complete if no other function can be added to it, preserving properties (1) (2) (3) (in particular prop. (3), orthogonality).

Theorem: if an orthogonal system has the Parseval Property (*), then it is complete.

Example:

$$f(x) \begin{cases} x & \text{for } x \in (-\pi; \pi) \\ 2\pi\text{-periodic} & \end{cases}$$

$$f(x) = \sum_{n=-\infty}^{+\infty} c_n e^{inx} \quad c_n = \frac{(-1)^n i}{n} \quad (n \neq 0)$$



Th. $\phi_n(x) = e^{inx}$ is complete in $L^2([- \pi; \pi])$ $n \in I \equiv \mathbb{Z}$ if

$$\sum_{\substack{n=-\infty \\ n \neq 0}}^{+\infty} \left| \frac{(-1)^n i}{n} \right|^2 \cdot \underbrace{\|\phi_n\|_2^2}_{2\pi} = \|f\|_2^2 \quad \downarrow \int_{-\pi}^{\pi} x^2 dx$$

$$2\pi \sum_{n \neq 0} \frac{1}{n^2} = \left[\frac{x^3}{3} \right]_{-\pi}^{\pi}$$

$$2\pi \left(\sum_{n=1}^{\infty} \frac{1}{n^2} + \sum_{n=-1}^{-\infty} \frac{1}{n^2} \right) = \frac{1}{3} (\pi^3 + \pi^3)$$

$$4\pi \overbrace{\sum_{n=1}^{\infty} \frac{1}{n^2}}^{\frac{\pi^2}{6}} = \frac{2}{3} \pi^3$$

$$\frac{4\pi^3}{6} = \frac{2}{3} \pi^3$$

The equality is true,
so the orthogonal system is complete!

Complete Orthogonal Systems of functions in L^2 , $\phi_\kappa(x)$

We say that $\{\phi_\kappa\}$, $\kappa \in I$, is ortho-normal if

$$\|\phi_\kappa\|_2 = 1 \quad \forall \kappa \in I$$

N.B. Any orthogonal system can be normalized just multiplying $\phi_\kappa(x)$ by a constant $\frac{1}{\|\phi_\kappa\|_2}$

Th. 2.3.1 (p. 33) (minimizing property of Fourier Coeff's)

Suppose $\{\phi_\kappa(x)\}_{\kappa \in I}$, $x \in (a, b)$, is an orthogonal system, $f \in L^2(a, b)$, $f \sim \sum_{\kappa \in I} c_\kappa \phi_\kappa(x)$

$$\Rightarrow \|f(x) - \sum_{\kappa=-m}^n c_\kappa \phi_\kappa(x)\|_2 \stackrel{(*)}{\leq} \|f(x) - \sum_{\kappa=-m}^n d_\kappa \phi_\kappa(x)\|_2$$

and the $=$ in $(*)$ holds $\Leftrightarrow c_\kappa = d_\kappa \quad \forall \kappa$

N.B. c_κ is the (generalized) Fourier coefficient of $f(x)$ with respect to the orthogonal system $\{\phi_\kappa(x)\}$ i.e.

$$c_\kappa = \frac{1}{\|\phi_\kappa\|_2^2} \cdot (f, \phi_\kappa)$$

$(*)$ means the "square deviation" of $f(x)$ from

$\sum_{\kappa=-m}^n d_\kappa \phi_\kappa(x)$ is minimal when $d_\kappa = c_\kappa$, where

c_κ are the Fourier coeff. of f w.r.t. $\{\phi_\kappa\}$
(with respect to)

Proof: $\left\| f - \sum_{k=-n}^n d_k \phi_k \right\|_2^2 = \left(f - \sum_{k=-n}^n d_k \phi_k, f - \sum_{k=-n}^n d_k \phi_k \right) =$

$\overset{\text{II}}{f(x)} \quad \overset{\text{III}}{\phi_k(x)}$

(Scalar product)

using
distributive property = $(f, f) - (f, \sum_{k=-n}^n d_k \phi_k) - (\sum_{k=-n}^n d_k \phi_k, f) +$

 $+ (\sum_{k=-n}^n d_k \phi_k, \sum_{k=-n}^n d_k \phi_k) = \|f\|_2^2 - \sum_{k=-n}^n (f, d_k \phi_k) +$
 $- \sum_{k=-n}^n (d_k \phi_k, f) + \sum_{j=-n}^n \sum_{k=-n}^n (d_k \phi_k, d_j \phi_j) =$

Remark

$$(\alpha f, g) = \alpha (f, g)$$

$$(f, \alpha g) = \bar{\alpha} (f, g)$$

$$= \|f\|_2^2 - \sum_{k=-n}^n \bar{d}_k (f, \phi_k) - \sum_{k=-n}^n d_k (\phi_k, f) +$$
 $+ \sum \sum d_k \bar{d}_j (\phi_k, \phi_j) = (*) \quad (\phi_k, \phi_j) \neq 0 \quad \text{when } k=j$

Remark

$$(f, \phi_k) = c_k \| \phi_k \|_2^2 \quad (\text{Fourier coeff. } c_k)$$

$$(\phi_k, f) = \overline{(f, \phi_k)} = \bar{c}_k \| \phi_k \|_2^2$$

$$(*) = \|f\|_2^2 - \sum_{k=-n}^n \bar{d}_k c_k \| \phi_k \|_2^2 - \sum_{k=-n}^n d_k \bar{c}_k \| \phi_k \|_2^2 +$$
 $+ \sum_{k=-n}^n d_k \bar{d}_k \| \phi_k \|_2^2 = + \sum_{k=-n}^n c_k \bar{c}_k \| \phi \|_2^2 - \sum_{k=-n}^n c_k \bar{c}_k \| \phi \|_2^2$

collect these terms

 $= \|f\|_2^2 - \sum_{k=-n}^n c_k \bar{c}_k \| \phi_k \|_2^2 + \sum_{k=-n}^n (c_k - d_k)(\bar{c}_k - \bar{d}_k) \| \phi_k \|_2^2 =$
 $= \|f\|_2^2 - \sum_{k=-n}^n |c_k|^2 \| \phi_k \|_2^2 + \sum_{k=-n}^n |c_k - d_k|^2 \| \phi_k \|_2^2 \quad (\star)$

This quantity is minimal $\iff c_k = d_k \text{ for } k = -n, -n+1, \dots, n-1, n$

(QED)

Furthermore If $c_n = d_n$ $S_m f(x) = \sum_{n=-m}^m c_n \phi_n(x)$

Partial Sums Order n (bilateral)
of our F.S. (generalized)

Then the last term in $(*)$ formula is $= 0 \Rightarrow$

$$\|f(x) - S_m f(x)\|_2^2 = \|f\|_2^2 - \sum_{n=-m}^m |c_n|^2 \|\phi_n\|_2^2 \quad (\text{proof of corollary})$$

This observation ("furthermore") implies proof of Bessel's

Inequality: $\|f\|_2^2 \geq \sum_{n=-m}^m |c_n|^2 \|\phi_n\|_2^2$

Def. we say that $\{\phi_n\}_{n \in I}$ has the Parseval Property

if holds in $\|f\|_2^2 = \sum_{n \in I} |c_n|^2 \|\phi_n\|_2^2$

Theorem (no proof in this 5 credits course!)

An orthogonal system of functions $\{\phi_n(x)\}_{n \in I}$ is complete \Leftrightarrow it has the Parseval Property.

Let's revisit these facts in the classical case when

$$\{\phi_n(x) = e^{inx}\} \text{ and } I = \mathbb{Z} = \{-3, -2, -1, 0, 1, 2, 3, \dots\}$$

Also, we observed that, in the classical case, we can write down in 2 forms the F.S. of $f(x)$

$$v \sum_{n=-\infty}^{+\infty} c_n e^{inx} = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

N.B. $\{\cos nx, \sin nx, \frac{1}{2}\}_{n=1,2,3,\dots}$ is a complete

orthogonal system in $L^2(-\pi; \pi)$. And, of course,

also $\{e^{inx}\}_{n \in \mathbb{Z}}$.

Parseval $\|f\|_2^2 = \int_{-\pi}^{\pi} |f(x)|^2 dx = \sum_{n=-\infty}^{+\infty} |c_n|^2 \|\phi_n\|_2^2$

$$\{e^{inx}\}$$

$$2\pi \forall n$$

$$\text{i.e. } \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(x)|^2 dx = \sum_{n=-\infty}^{+\infty} |c_n|^2$$

Parseval $\{\cos nx, \sin nx, \frac{1}{2}\}$

We proved that

$$b_0 = 0$$

$$c_n = \frac{a_n - i b_n}{2} \quad \text{for } n > 0 \quad (c_0 = a_0/2)$$

$$c_n = \frac{a_n + i b_n}{2} \quad \text{for } n < 0$$

$$a_n = c_n + \bar{c}_n \quad \text{for } n > 0 \quad (a_0 = 2c_0)$$

$$b_n = i(c_n - \bar{c}_n) \quad \text{for } n > 0 \quad (b_0 = 0)$$

$$\sum_{n=-\infty}^{+\infty} |c_n|^2 = \sum_{n=1}^{\infty} |\bar{c}_n|^2 + |c_0|^2 + \sum_{n=1}^{\infty} |c_n|^2 = \text{we substitute the previous relationships} =$$

$$= \sum_{n=1}^{\infty} \frac{|a_n + i b_n|^2}{4} + \frac{|a_0|^2}{4} + \sum_{n=1}^{\infty} \frac{|a_n - i b_n|^2}{4} =$$

by def. of
complex number
modulus:

$$= \frac{|a_0|^2}{4} + \frac{1}{2} \sum (|a_n|^2 + |b_n|^2), \text{ because } \frac{1}{4} [(a_n + i b_n)(\bar{a}_n - i \bar{b}_n)$$

$$|\bar{z}| = \sqrt{x^2 + y^2}$$

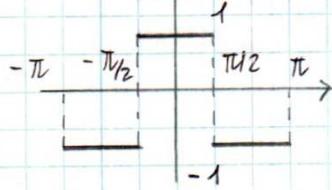
$$+ (a_n - i b_n)(\bar{a}_n + i \bar{b}_n)] =$$

$$= \frac{1}{4} [|a_n|^2 - i a_n \bar{b}_n + i b_n \bar{a}_n + |b_n|^2 + |a_n|^2 + i a_n \bar{b}_n + \\ - i b_n \bar{a}_n + |b_n|^2] = \frac{1}{2} (|a_n|^2 + |b_n|^2)$$

$$\boxed{(\text{Real}) \text{ Parseval} \quad \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(x)|^2 dx = \frac{|a_0|^2}{4} + \frac{1}{2} \sum_{n=1}^{\infty} (|a_n|^2 + |b_n|^2)}$$

example

$$\text{F.S. } f(x) = \frac{4}{\pi} \left(\cos x - \frac{\cos 3x}{3} + \frac{\cos 5x}{5} - \dots \right)$$



Square Wave

Parseval
 $(b_n = 0)$
 $(a_0 = 0)$

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} |f(x)|^2 dx =$$

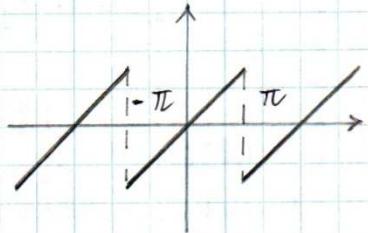
$$= \frac{1}{2} \frac{16}{\pi^2} \left(\frac{1}{1^2} + \frac{1}{3^2} + \frac{1}{5^2} + \frac{1}{7^2} + \dots \right)$$

$$\underbrace{1}_{\text{where } a_n = \frac{(-1)^{n+1}}{\pi}} \cdot \frac{4}{\pi}$$

$$\underbrace{\frac{1}{2\pi} \int_{-\pi}^{\pi} 1 dx}_{\text{1}} \Rightarrow \frac{1}{1^2} + \frac{1}{3^2} + \frac{1}{5^2} + \dots = \frac{\pi^2}{8}$$

Other ex. sawtooth

$$f(x) = 2 \left(\sin x - \frac{\sin 2x}{2} + \frac{\sin 3x}{3} - \dots \right)$$



$$\frac{1}{2\pi} \int_{-\pi}^{\pi} x^2 dx = \frac{1}{2} \sum_{n=1}^{\infty} 2^2 \frac{1}{n^2}$$

$$\frac{1}{2\pi} \left[\frac{x^3}{3} \right]_{-\pi}^{\pi} = \frac{4}{2} \left(\frac{1}{1^2} + \frac{1}{2^2} + \frac{1}{3^2} + \dots \right)$$

$$\frac{1}{4\pi} \left(\frac{\pi^3 - (-\pi)^3}{3} \right) = \frac{1}{1^2} + \frac{1}{2^2} + \frac{1}{3^2} + \dots$$

$$\Rightarrow \frac{\pi^2}{6} = \frac{1}{1^2} + \frac{1}{2^2} + \frac{1}{3^2} + \dots$$

Theorem

Schwarz Inequality

(Cauchy
Schwarz
Bunjakowski)

$$|(f, g)| \leq \|f\|_2 \|g\|_2 \quad (\text{scalar product})$$

$$\text{i.e. } \left| \int_a^b f(x) \overline{g(x)} dx \right| \leq \left(\int_a^b |f(x)|^2 dx \right)^{1/2} \left(\int_a^b |g(x)|^2 dx \right)^{1/2}$$

N.B. this inequality is "obvious" if we take for granted the euclidean (∞ dimensional) nature of scalar products and L^2 norms, otherwise not so obvious...

Remark The Schwarz Inequality is also true if our scalar product contains a weight function $w(x) > 0$

$$\text{i.e. } (f, g) \stackrel{\text{def.}}{=} \int_a^b f(x) \overline{g(x)} w(x) dx$$

in particular $w(x) = 1$

$$\left| \int_a^b f(x) \overline{g(x)} w(x) dx \right| \leq \left(\int_a^b |f(x)|^2 w(x) dx \right)^{1/2} \left(\int_a^b |g(x)|^2 w(x) dx \right)^{1/2}$$

Proof: set $\vartheta = \operatorname{Arg}(f, g)$

[N.B. if (f, g) is real $\vartheta = 0$ if $(f, g) > 0$
 $\vartheta = \pi$ if $(f, g) < 0$]

Then $e^{-i\vartheta} (f, g) = (e^{-i\vartheta} f, g) \equiv |(f, g)|$

produces a rotation ϑ coincides with the abs. val.

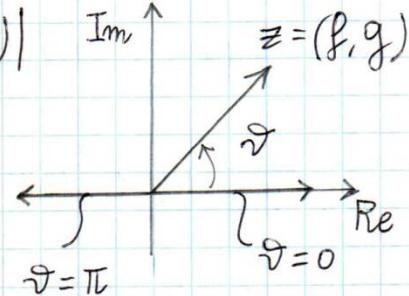
in this case $\vartheta = 0, \pi$

$$z_1 = s_1 e^{i\vartheta_1}$$

$$z_1 z_2 = s_1 s_2 e^{i(\vartheta_1 + \vartheta_2)}$$

$$z_2 = s_2 e^{i\vartheta_2}$$

DE MOIVRE FORMULA



Now, $e^{-i\vartheta}$ is a complex number of argument $-\vartheta$ and $|e^{-i\vartheta}| = 1$

So the effect of multiplying $z \in \mathbb{C}$ by $e^{-i\vartheta}$ is to rotate the vector z (in the complex plane) of an angle $-\vartheta$

If $\operatorname{Arg} z = \vartheta$, then $e^{-i\vartheta} z = |z|$

Also, in this case, ($\vartheta = 0, \pi$)

$$(f e^{-i\vartheta}, g) = \operatorname{Re}(f e^{-i\vartheta}, g) = |(f, g)|$$

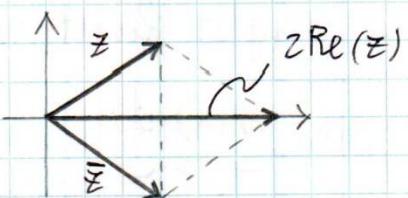
Now we go on with our proof:

$$0 \leq \|f e^{-i\vartheta} + \lambda g\|_2^2 \quad \begin{aligned} &\text{(Where } \lambda \in \mathbb{R} \text{)} \\ &\vartheta = \operatorname{Arg}(f, g) \end{aligned}$$

$$= (f e^{-i\vartheta} + \lambda g, f e^{-i\vartheta} + \lambda g) = \text{distributive property} =$$

$$= (f e^{-i\vartheta}, f e^{-i\vartheta}) + \lambda (f e^{-i\vartheta}, g) + \lambda (g, f e^{-i\vartheta}) + \lambda^2 (g, g) =$$

$$= \|f e^{-i\vartheta}\|_2^2 + \lambda \operatorname{Re}(f e^{-i\vartheta}, g) + \lambda^2 \|g\|_2^2 =$$



$$= \|f\|_2^2 + 2\lambda |(f, g)| + \lambda^2 \|g\|_2^2$$

\leq polynomial of degree 2 in the variable λ

(N.B. coeff. of λ^2 is $\|g\|_2^2 > 0$) $\Rightarrow \Delta \leq 0$

$$\frac{\Delta}{4} = \frac{b^2 - 4ac}{4} = \left(\frac{b}{2}\right)^2 - ac$$

$$\Rightarrow |(f, g)|^2 - \|f\|_2^2 \|g\|_2^2 \leq 0$$

$$|(f, g)|^2 \leq \|f\|_2^2 \|g\|_2^2$$

$$|(f, g)| \leq \|f\|_2 \|g\|_2 \quad (\text{QED})$$

$$|(f, g)| \leq \|f\|_2 \|g\|_2$$

True also on $L^2(\mathbb{R})$ or $L^2([0, +\infty))$,

so this inequality is relevant also
for the Fourier Transform and
Laplace Transform.

Lemma 2.4.2

(P.36)

$$(f, g) = \int_a^b f(x) \overline{g(x)} w(x) dx = \sum_{n \in I} c_n \overline{d_n} \|\phi_n\|_2^2$$

where c_n and d_n are the Fourier coeff. of f and g
with respect to the orthogonal system $\{\phi_n(x)\}$

$$f \sim \sum_{n \in I} c_n \phi_n(x) \quad \text{with } c_n = \frac{1}{\|\phi_n\|_2^2} (f, \phi_n)$$

$$\text{and } g \sim \sum_{n \in I} d_n \phi_n(x)$$

$$d_n = \dots$$

This formula (Parseval-Plancherel) in the special case $f = g$

gives $\int_a^b |f(x)|^2 dx = \sum_{n \in \mathbb{Z}} |c_n|^2 \|\phi_n\|_2^2$

Proof $|(\bar{f} - S_N f, g)| \leq \|f - S_N f\|_2 \|g\|_2 \rightarrow 0$

$$S_N f = \sum_{n=-N}^N c_n \phi_n(x)$$

(because $\|g\|_2 > 0$ and

$$\lim_{N \rightarrow \infty} \|f - S_N f\|_2 = 0$$

$$(f, g) = (\bar{f} - S_N f + S_N f, g) = (\bar{f} - S_N f, g) + (S_N f, g)$$

therefore

$$(f, g) = \lim_{N \rightarrow \infty} (S_N f, g) = \lim_{N \rightarrow \infty} \left(\sum_{n=-N}^N c_n \phi_n(x), g(x) \right) =$$

$$= \lim_{N \rightarrow \infty} \sum_{n=-N}^N c_n \underbrace{(\phi_n, g)}_{\parallel \cdot \parallel} = \lim_{N \rightarrow \infty} \sum_{n=-N}^N c_n \overline{d_n} \underbrace{\|\phi_n\|_2^2}_{\text{FROM THE DEFINITION OF THE GENERALIZED COEFF'S}} = \\ = \sum_{n=-\infty}^{+\infty} c_n \overline{d_n} \|\phi_n\|_2^2$$

(QED)

Th. (No proof, but easy consequence of

what we have done today) if $f \in L^1(a, b)$

and its F.S. is $f \sim \sum_{n \in \mathbb{Z}} c_n \phi_n(x)$ suppose $[a, b] \subset [a, b]$

then

$$\int_a^b f(x) dx = \int_a^b \sum_{n \in \mathbb{Z}} c_n \phi_n(x) dx = \sum_{n \in \mathbb{Z}} c_n \int_a^b \phi_n(x) dx$$

{if $[a, b] \subset [a, b]$ finite intervals and $f \in L^1([a, b]) \cap L^2([a, b])$ }
 {then it's ok to exchange $\int_a^b \dots$ with F. Series}

Remark In applications to ODE and PDE we would

also like to use $[f(x)]' = [\sum_{n \in \mathbb{Z}} c_n \phi_n(x)]' = \sum_{n \in \mathbb{Z}} c_n \phi_n'(x)$

In general this is more delicate than exchanging F.S. and integrals... But it's OK if $c_k \rightarrow 0$ "quickly enough" (equivalent to $f(x)$ being "smooth enough").

Some families of orthogonal systems $\{\phi_\kappa(x)\}_{\kappa \in I}$

which are not the classical $\{e^{ikx}\}_{k \in \mathbb{Z}}$ or the equivalent $\{\cos kx, \sin kx, \frac{1}{2}\}_{k=1,2,3,\dots}$

Orthogonal Polynomials (pg. 45-59)

Idea: on a finite interval $[a; b]$ the power functions $1, x, x^2, x^3, \dots$ are linearly independent, i.e.,

$$\sum_{n=0}^N a_n x^n = 0 \text{ for all } x \in [a; b]$$

$$\Leftrightarrow a_n = 0 \text{ for } n = 0, 1, 2, \dots, N$$

This fact is somehow connected to the fact that it is possible to approximate large spaces of functions using polynomials. For example if $f \in L^p([a; b])$ $\forall \epsilon > 0 \exists n \in \mathbb{N}$ and a polynomial $p(x)$ of degree n such that

$$\|f(x) - p(x)\|_{L^p([a; b])} < \epsilon \quad \left\{ \begin{array}{l} \text{N.B. the } L^\infty \text{ case is} \\ \text{UNIFORM approximation} \end{array} \right\}$$

Given a scalar product on $L^2([a; b], w)$

$$(f, g) \stackrel{\text{def}}{=} \int_a^b f(x) \overline{g(x)} w(x) dx$$

usually is not true that the power functions are orthogonal with respect to each other.

BUT it is possible to find a sequence of polynomials

$p_0(x), p_1(x), p_2(x), p_3(x), \dots$, where each $p_n(x)$

is a polynomial of (exact) degree n (the coeff. of x^n in $p_n(x)$ is $\neq 0$) such that $(p_n(x), p_j(x)) = \int_a^b p_n(x) \overline{p_j(x)} w(x) dx =$

$$\text{Graph of } p_n(x) \quad \text{Graph of } p_j(x) \quad = \begin{cases} 0 & \text{if } n \neq j \\ \alpha_{n>0} & \text{if } n = j \end{cases}$$

Furthermore if we fix $[a; b]$ and a weight $w(x) > 0$ for $x \in [a; b]$ (possibly $w(x) = 1$), the sequence $p_0(x), p_1(x), p_2(x), \dots$ is essentially UNIQUE.

"essentially" in the sense that each $p_n(x)$ can be dilated changing it into $B_n p_n(x)$

If we normalize, e.g. choosing $\|p_n\|_2 = 1 \quad \forall n = 0, 1, 2, \dots$

then the sequence is unique. (other normalizations are possible, e.g., $p_n(0) = 1$)

Problem Given $[a; b]$ and $w(x)$, how do we compute $p_0(x), p_1(x), p_2(x), \dots$ orthogonal polynomials with respect to $(f, g) = \int_a^b f(x) \overline{g(x)} w(x) dx$?

Several answers:

1st: "Bulldozer" answer: Gram-Schmidt Orthogonalization

Advantage: it "always" works (for any given $[a; b] w(x)$) starting from the non-orthogonal but independent system $1, x, x^2, x^3, \dots$

Disadvantage: boring, inefficient, slow (even on computers)

2nd possibility: It can be proven that all orthogonal polynomials satisfy a 3 terms recursive formula:

$$P_m = A_m P_{m-1}^{(x)} + B_m P_{m-2}^{(x)}$$

Recipe: compute $P_0(x)$ and $P_1(x)$. From P_0 and $P_1(x)$ compute $P_2(x)$ using *

From P_1 and P_2 compute P_3 etc. ...

3rd method: Direct Formulae given n produce P_n

4th method: Semidirect Formulae (Rodriguez)

$$P_n(x) = \frac{d^n}{dx^n} (\text{something})$$

Other methods...

Depending on $[a; b]$ and $w(x)$ the "smart" method could change.

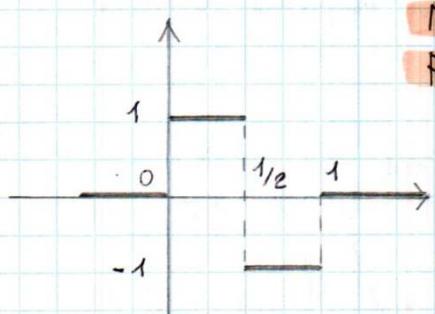
In this course we'll see only a small sample of these methods (mostly following the book)

Wavelets (in particular Haar Basis) [ALSO p.62]

Theorem (Haar)

$$\psi(x) = \begin{cases} 1 & \text{if } x \in (0, \frac{1}{2}) \\ -1 & \text{if } x \in (\frac{1}{2}, 1) \\ 0 & \text{elsewhere} \end{cases}$$

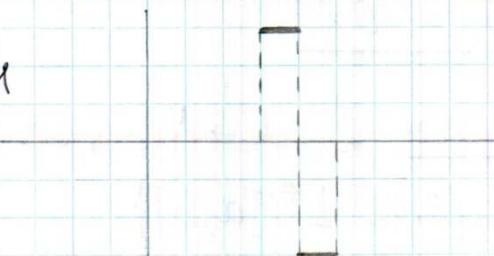
MOTHER
FUNCTION



(ψ is "one tooth" of a square wave)

Def. $\psi_{jn}(x) = 2^{j/2} \psi(2^j x - n)$ for $j \in \mathbb{Z} \ n \in \mathbb{Z}$

ψ_{jn}



$\psi_{jn}(x)$ is a complete orthonormal system in $L^2(\mathbb{R})$

in other words any $f \in L^2(\mathbb{R})$ can be written

$$f(x) = \sum_{j,k} c_{jk} \psi_{jk}(x), \text{ where the "Haar coefficients"}$$

(Generalized F.S.)

$$c_{jk} = (f, \psi_{jk}) =$$

(N.B. $\|\psi_{jk}\|_2^2 = 1$
because $\{\psi_{jk}\}$ is
an orthonormal system)

$$= \int_{-\infty}^{+\infty} f(x) \overline{\psi_{jk}(x)} dx = \int_{\text{interval support}} f(x) \psi_{jk}(x) dx$$

(real function) of ψ_{jk}

Advantage of Haar system: very simple $\psi(x)$

Disadvantages: many coefficients are needed to get a good approximation of $f(x)$. Properties of $f(x)$ (like smoothness) are mapped into c_{jk} in a complicated way.

Modern approach: choose better $\psi(x)$ mother functions such that $\psi_{jk}(x) = 2^{j/2} \psi(2^j x - k)$ for $j, k \in \mathbb{Z}$ is a complete orthonormal system in $L^2(\mathbb{R})$.

Other big families of orthogonal systems are given by eigenfunctions of specific operators.

COMPARE
PAGE
41R &
44

Def. If A is an $n \times n$ matrix it defines linear operator: $\mathbb{R}^n \rightarrow \mathbb{R}^n$

$$\begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{pmatrix} = \begin{pmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ \vdots & \vdots & & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nn} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ \vdots \\ x_n \end{pmatrix}$$

α is an EIGENVALUE and \vec{x} is an EIGENVECTOR of this operator if $A\vec{x} = \alpha\vec{x}$

Def. Given a linear operator $L: A \rightarrow A$, where A is a space of functions (e.g. $L^p([a,b])$...)

Linear

$$\mathcal{L}(\alpha f(x) + \beta g(x)) = \alpha \mathcal{L}(f(x)) + \beta \mathcal{L}(g(x))$$

(ex. of linear operators: derivative, n-th derivative, Fourier transform, Laplace Transform ...)

If $\mathcal{L}(f(x)) = \alpha \cdot f(x)$, then we say that α is an EIGENVALUE and f is an EIGENFUNCTION of \mathcal{L} .

Ex. $f(x) = e^{\alpha x}$ $\mathcal{L} = \text{derivative}$

$$\mathcal{L}(e^{\alpha x}) = \alpha e^{\alpha x} \quad \text{eigenvalue} = \alpha$$

$\mathcal{L} = \text{second derivative}$ $f(x) = \cos x, \sin x$
eigenfunctions

$$(\sin x)'' = -\sin x \quad \mathcal{L}(\sin x) = (-1)\sin x \\ \text{eigenvalue} = -1$$

$$\mathcal{L}(\cos x) = -\cos x \quad \text{eigenvalue} = -1$$

N.B. $\mathcal{L}(3\cos x) = 3(-\cos x) = (-1)(3\cos x)$

↑ same eigenvalue

Ex. $\mathcal{L} = \text{Fourier Transform}$

$$\mathcal{L}(e^{-ax^2}) = \sqrt{\frac{\pi}{a}} e^{-t^2/4a} \quad (\text{Gaussian function})$$

$$a > 0 \quad \text{let's choose } a = \frac{1}{2}$$

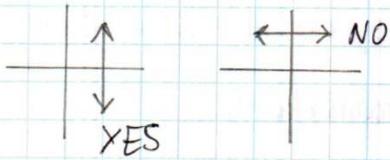
$$\mathcal{L}(e^{-\frac{1}{2}x^2}) = \sqrt{2\pi} e^{-\frac{1}{2}t^2}$$

$$\text{so } f(x) = e^{-x^2/2} \text{ eigenfunction eigenvalue} = \sqrt{2\pi}$$

Rem.1: we have to choose $a = \frac{1}{2}$ to have $f(x)$ eigenfunction

of the F.T.; the corresponding eigenvalue is $\sqrt{2\pi}$ BUT =

$$\mathcal{L}(e^{-izx}) = c \sqrt{2\pi} e^{-iz^2}$$



Rem. 2: (more details later) it is possible to find polynomials
P.43R $h_0(x), h_1(x), h_2(x), \dots, h_n(x)$ of degree n such that

$$\text{F.T. } (h_n(x) e^{-x^2/2}) = \alpha_n \underbrace{h_n(x)}_{\text{HERMITE POLYNOMIALS}} \underbrace{e^{-x^2/2}}_{\text{HERMITE FUNCTIONS}}$$

It is possible to prove also that these eigenfunctions of the F.T. $h_n(x) e^{-x^2/2}$ are a complete orthogonal system in $L^2(\mathbb{R})$.



Orthogonal polynomials

$p_m(x) = c_n x^n + \dots + c_1 x + c_0$ (linear combination of powers
degree = n = highest power)

Given $[a; b]$ (also $[0; +\infty)$, also $(-\infty; +\infty)$) and a weight (i.e. a function $w(x) > 0$ for $x \in [a; b]$) we have a scalar product $(f, g) = \int_a^b f(x) \overline{g(x)} w(x) dx$; we also have an "essentially" unique sequence of polynomials of degree $n = 0, 1, 2, 3, \dots$ such that it is an orthogonal system,

$$\text{i.e. } (p_m, p_m) = \int_a^b p_m(x) \overline{p_m(x)} w(x) dx = \begin{cases} 0 & \text{if } m \neq m \\ \|p_m\|_2^2 & \text{if } m = m \end{cases}$$

Depending on the weight choice we get different solutions.

case n. 1 Legendre Polynomials: $[a; b] = [-1; 1]$ $w(x) = 1 \forall x \in [-1, 1]$
our scalar product is $(f, g) = \int_{-1}^1 f(x) \overline{g(x)} dx$. They are:

$$p_0(x) = 1; \quad p_1(x) = x; \quad p_2(x) = \frac{3}{2}x^2 - \frac{1}{2}; \quad p_3(x) = \frac{5}{2}x^3 - \frac{3}{2}x;$$

$$p_4(x) = \frac{35}{8}x^4 - \frac{15}{4}x^2 + \frac{3}{8}$$

(are chosen in such a way that they make oscillation in $[-1; 1]$)

Let's check orthogonality in one case $(p_2, p_3) = \int_{-1}^1 \left(\frac{3}{2}x^2 - \frac{1}{2}\right) \left(\frac{5}{2}x^3 - \frac{3}{2}x\right) dx$

$$= \int_{-1}^1 \underbrace{\frac{15}{4}x^5 - \frac{9}{4}x^3 - \frac{5}{4}x^3 + \frac{3}{4}x}_{-\frac{14}{4}x^3} dx = \left[\frac{15}{4} \frac{x^6}{6} - \frac{14}{4} \frac{x^4}{4} + \frac{3}{4} \frac{x^2}{2} \right]_{-1}^1 =$$

$$= \frac{5}{8} - \frac{7}{8} + \frac{3}{8} - (\text{same things}) = 0$$

$$(P_2, P_2) = \int_{-1}^1 \left(\frac{3}{2}x^2 - \frac{1}{2}\right)^2 dx > 0$$

These are polynomials, but, in the interval they act, they are imitating sine & cosine (oscillation, cancellation effect...)

See figure 3.1, page 47 of the book.

To obtain O.P. we could apply the Gram-Schmidt orthogonalization to the sequence of powers $1, x, x^2, x^3, \dots$

(N.B. these polynomials are linearly independent but they are not orthogonal w.r.t. $(f, g) = \int_{-1}^1 f(x) \overline{g(x)} dx$)

There are more efficient methods to obtain the same sequence of polynomials.

It can be proven that
$$P_m(x) = \frac{1}{2^m \cdot m!} \frac{d^m}{dx^m} (x^2 - 1)^m \quad (*)$$
 (Rodriguez)

(hidden recursion...)

↳ because you can't compute directly the n -th derivative.

$P_0(x) = 1$ From the formula (*) (Rodriguez):

$$P_1(x) = \frac{1}{2^1 \cdot 1!} \frac{d}{dx} (x^2 - 1) = \frac{1}{2} \cdot 2x = x \quad \checkmark$$

$$P_2(x) = \frac{1}{2^2 \cdot 2!} \frac{d^2}{dx^2} (x^2 - 1)^2 = \frac{1}{8} \frac{d^2}{dx^2} (x^4 - 2x^2 + 1) = \frac{1}{8} (4x^3 - 4x)' =$$

$$= \frac{1}{8} (12x^2 - 4) = \frac{3}{2}x^2 - \frac{1}{2} \quad \checkmark$$

Theorem (no proof in this course)

$$(\star) \quad [n P_n(x) - x(2n-1) P_{n-1}(x) + (n-1) P_{n-2}(x)] = 0 \quad \begin{matrix} \text{(3-terms)} \\ \text{recursion} \end{matrix}$$

Special case of Theorem*: Almost all systems of O.P. satisfy a 3 term recursion (0|||^o)

From (\star) if we know the first two cases $P_0(x) = 1$, $P_1(x) = x$ then recursively we can compute $P_2(x)$, $P_3(x)$, ...

Legendre Polynomials

Rodriguez $P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} (x^2 - 1)^n$

$$P_0(x) = 1$$

$$P_1(x) = \frac{1}{2} \frac{d}{dx} (x^2 - 1) = x$$

$$P_2(x) = \frac{1}{2^2 2!} \frac{d^2}{dx^2} (x^2 - 1)^2 = \frac{1}{8} \frac{d}{dx} [2(x^2 - 1) \cdot 2x] =$$

$$= \frac{1}{8} [4x^3 - 4x] = \frac{1}{8} [12x^2 - 4] = \frac{3}{2}x^2 - \frac{1}{2}$$

$$P_3(x) = \frac{1}{2^3 3!} \frac{d^3}{dx^3} (x^2 - 1)^3 = \frac{1}{8 \cdot 6} \frac{d^2}{dx^2} [3(x^2 - 1)^2 \cdot 2x] =$$

$$(x^4 + 1 - 2x^2) 6x$$

$$= \frac{1}{48} \frac{d}{dx} (30x^4 + 6 - 36x^2) = (6x^5 + 6x - 12x^3)$$

$$= \frac{1}{48} (120x^3 - 72x) = \frac{5}{2}x^3 - \frac{3}{2}x$$

$$P_4(x) = \frac{1}{2^4 4!} \frac{d^4}{dx^4} (x^2 - 1)^4 = \frac{1}{384} \frac{d^3}{dx^3} [8x(x^2 - 1)^3] =$$

$$= \frac{1}{384} \frac{d^2}{dx^2} [8(x^2 - 1)^3 + 8x \cdot 3(x^2 - 1)^2 \cdot 2x] =$$

$$48x^2(x^2 - 1)^2$$

$$= \frac{1}{384} \frac{d}{dx} [24(x^2 - 1)^2 \cdot 2x + 96x(x^2 - 1)^2 + 48x^2 \cdot 2(x^2 - 1) \cdot 2x] =$$

$$144x(x^2 - 1)^2 + 192x^3(x^2 - 1)$$

$$= \frac{1}{384} \cdot [144(x^2 - 1)^2 + 144x \cdot 2(x^2 - 1) \cdot 2x + 576x^2(x^2 - 1) +$$

$$+ 192x^3 \cdot 2x] =$$

$$= \frac{1}{384} [144x^4 + 144 - 288x^2 + 576x^4 - 576x^2 + 576x^4 - 576x^2 + 384x^4]$$

$$\dots - 15\sqrt{2}x^4 - 11\sqrt{2}x^2 - 11\sqrt{2} - 35x^4 - 15\sqrt{2}x^2 + 384x^4$$

3 term recursion

$$n P_n(x) - x(2n-1)P_{n-1}(x) + (n-1)P_{n-2}(x) = 0$$

$$P_0(x) = 1 \quad P_1(x) = x$$

$$2P_2 - x(2 \cdot 2 - 1)x + (2 - 1)1 = 0$$

$$2P_2 = 3x^2 - 1 \quad P_2 = \frac{3x^2}{2} - \frac{1}{2}$$

$$3P_3 - x(6 - 1)\left(\frac{3}{2}x^2 - \frac{1}{2}\right) + (3 - 1)x = 0$$

$$3P_3 = \frac{15}{2}x^3 - \frac{5}{2}x - 2x = \frac{15}{2}x^3 - \frac{9}{2}x$$

$$P_3 = \frac{5}{2}x^3 - \frac{3}{2}x$$

$$4P_4 - x(8 - 1)\left(\frac{5}{2}x^3 - \frac{3}{2}x\right) + 3\left(\frac{3x^2}{2} - \frac{1}{2}\right) = 0$$

$$4P_4 - \frac{35}{2}x^4 + \frac{21}{2}x^2 + \frac{9}{2}x^2 - \frac{3}{2} = 0$$

$$4P_4 = \frac{35}{2}x^4 - 15x^2 + \frac{3}{2}$$

$$P_4 = \frac{35}{8}x^4 - \frac{15}{4}x^2 + \frac{3}{8}$$

HIGHLY MORE EFFECTIVE PROCEDURE

Exercise: compute the first 5 or 6 (or 7) Legendre Polynomials both with Rodriguez Formula and the (*) recursion comparing efficiency.

A little detour about recursive method: for example, Fibonacci Numbers $c_0, c_1, c_2, c_3, \dots$

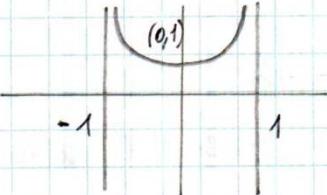
$$c_0 = c_1 = 1 \text{ and } c_n = c_{n-1} + c_{n-2}$$

$$c_2 = 1+1=2 \quad c_3 = 2+1=3 \quad c_4 = \dots$$

exercise: either on your own (better) or cheating (Google) find a direct closed formula for the n -th Fibonacci Number c_n (hint: it is the linear combination of two exponentials one increasing to $\rightarrow \infty$ and the other one decreasing to 0) (hint: Golden Ratio)

Sometimes (but not always) it is possible, like in the case of Fibonacci, to obtain a "closed formula" from the recursive definition.

Case 2: Chebycheff Polynomials $T_n(x)$ orthogonal system of polynomials in $(-1; 1)$ with the weight

$$W(x) = \frac{1}{\sqrt{1-x^2}}$$


$$(f, g) = \int_{-1}^1 f(x) \overline{g(x)} \frac{1}{\sqrt{1-x^2}} dx \quad (**)$$

$$T_0(x) = 1 \quad T_1(x) = x \quad T_2(x) = 2x^2 - 1$$

$$T_3(x) = 4x^3 - 3x \quad T_4(x) = 8x^4 - 8x^2 + 1 \quad T_5(x) = \dots$$

First method to compute these polynomials is Gram-Schmidt apply to $1, x, x^2, x^3, \dots$ w.r.t. (**)

2nd method

$$T_{n+1}(x) - 2xT_n(x) + T_{n-1}(x) = 0$$

(try it)

$$T_2(x) - 2xT_1(x) + T_0(x) = 0$$

Theorem "closed formula" for Chebycheff Polynomials

$$\begin{aligned} T_n(x) &= \cos(n \arccos x) = \dots = \\ &= x^n - \binom{n}{2} x^{n-2} (1-x^2) + \binom{n}{4} x^{n-4} (1-x^2)^2 - \dots \end{aligned}$$

Case 3 : Hermite Polynomials (and Hermite Functions)

consider the interval $(-\infty; +\infty) \equiv \mathbb{R}$ $W(x) = e^{-x^2}$

$$(f, g) = \int_{-\infty}^{+\infty} f(x) \overline{g(x)} e^{-x^2} dx$$



We get Hermite Polynomials $H_n(x)$

$$H_{n+1}(x) - 2xH_n(x) + 2nH_{n-1}(x) = 0 \quad H_0(x) = 1 \quad H_1(x) = 2x$$

Rodriguez :
$$H_n(x) = (-1)^n e^{x^2} \frac{d^n}{dx^n} e^{-x^2}$$

$$\text{e.g. } H_1(x) = (-1)^1 e^{x^2} \frac{d}{dx} (e^{-x^2}) = -e^{x^2}(-2x)e^{-x^2} = 2x$$

$$\begin{aligned} H_2(x) &= (-1)^2 e^{x^2} \frac{d}{dx} (-2x e^{-x^2}) = e^{x^2}(-2e^{-x^2} - 2x(-2x)e^{-x^2}) \\ &= -2 + 4x^2 \end{aligned}$$

Important Remark Once we have generated the sequence

$H_0(x), H_1(x), H_2(x), \dots$ of Hermite Polynomials we know that

$$\int_{-\infty}^{+\infty} H_m(x) H_n(x) e^{-x^2} dx = \begin{cases} 0 & m \neq n \\ \|H_m\|_2^2 > 0 & \text{if } m = m \end{cases}$$

We can also consider the $h_n(x) = H_n(x) e^{-x^2/2}$ called

Hermite Functions (N.B. these functions are not polynomials, but they are polynomials times the square root of a weight)

$$h_0(x) = e^{-x^2/2}; h_1(x) = 2x e^{-x^2/2}; h_2(x) = (4x^2 - 2) e^{-x^2/2}$$